

THE 4TH MOMENT OF THE VELOCITY DISTRIBUTION FOR AN ANISOTROPIC SPHERICALLY SYMMETRIC DISTRIBUTION.

FOR A NON-ROTATING SPHERICALLY SYMMETRIC SELF-GRAVITATING SYSTEM, THE VLASOV EQUATION CAN BE WRITTEN

$$\frac{Df}{dt} = v_r \frac{\partial f}{\partial r} + \left(\frac{v_\theta^2 + v_\phi^2}{r} - \frac{d\Phi}{dr} \right) \frac{\partial f}{\partial v_r} + \frac{1}{r} (v_\theta^2 \cot \theta - v_r v_\theta) \frac{\partial f}{\partial v_\theta} - \frac{1}{r} (v_\phi v_r + v_\phi v_\theta \cot \theta) \frac{\partial f}{\partial v_\phi} = 0 \quad (1)$$

WE'RE INTERESTED IN THE FOURTH MOMENT OF THE VELOCITY, AND SO MULTIPLY (1) BY $(v_r, \theta, \phi, v_r, \theta, \phi, v_r, \theta, \phi)$ AND INTEGRATE OVER VELOCITY SPACE. OF THE 10 POSSIBLE MOMENTS GENERATED IN THIS WAY, THE INTERESTING ONES ARE:

$$\int (1) v_r^3 d^3v \Rightarrow \frac{d}{dr} (\rho \overline{v_r^4}) + 3\rho \overline{v_r^2} \frac{d\Phi}{dr} - \frac{3\rho}{r} \overline{v_r^2 (v_\theta^2 + v_\phi^2)} + \frac{2\rho}{r} \overline{v_r^4} = 0 \quad (2)$$

$$\int (1) v_r v_\theta^2 d^3v \Rightarrow \frac{d}{dr} (\rho \overline{v_r^2 v_\theta^2}) + \rho \overline{v_\theta^2} \frac{d\Phi}{dr} - \frac{\rho}{r} (\overline{v_\theta^4} + \overline{v_\theta^2 v_\phi^2}) + \frac{4\rho}{r} \overline{v_\theta^2 v_r^2} = 0 \quad (3)$$

FOR A SPHERICALLY SYMMETRIC SYSTEM, THERE ARE ONLY 2 INDEPENDENT COMPONENTS OF VELOCITY, v_r AND $v_E = (v_\theta^2 + v_\phi^2)^{1/2}$. TO CONVERT THE TERMS WITH v_θ AND v_ϕ IN TO v_E , WE WRITE THE VELOCITY PHASE SPACE AT ANY POINT IN SPHERICAL POLAR COORDINATES:

$$\left. \begin{aligned} v_\theta &= v \sin \xi \sin \eta \\ v_\phi &= v \cos \xi \sin \eta \\ v_r &= v \cos \eta \end{aligned} \right\} \Rightarrow v_E = v \sin \eta \quad \left. \begin{aligned} & \\ & \\ & \end{aligned} \right\} d^3v = v^2 \sin \eta d\eta d\xi dv$$

$$\begin{aligned} \Rightarrow \rho \overline{v_\theta^4} &= \iiint v^4 \sin^4 \xi \sin^4 \eta f \cdot v^2 \sin \eta d\eta d\xi dv \\ &= \int_{\xi=0}^{2\pi} \frac{1}{2\pi} \sin^4 \xi d\xi \cdot 2\pi \iint v^4 \sin^4 \eta f v^2 \sin \eta d\eta dv \\ &= \frac{3}{8} \cdot \rho \overline{v_E^4} \end{aligned}$$

AND SIMILARLY, $\rho \overline{v_\phi^4} = \frac{3}{8} \rho \overline{v_E^4}$; $\rho \overline{v_\theta^2 v_\phi^2} = \frac{1}{8} \rho \overline{v_E^4}$

$$\overline{v_\theta^2} = \overline{v_\phi^2} = \frac{1}{2} \overline{v_E^2}; \quad \overline{v_r^2 v_\theta^2} = \overline{v_r^2 v_\phi^2} = \frac{1}{2} \overline{v_r^2 v_E^2}$$

SUBSTITUTING THIS INFORMATION INTO EQUATIONS (2) + (3)

$$\Rightarrow \frac{d}{dr} (\rho \bar{v}_r^4) + 3\rho \bar{v}_r^2 \frac{d\Phi}{dr} - \frac{3\rho}{r} \bar{v}_r^2 \bar{v}_E^2 + \frac{2\rho}{r} \bar{v}_r^4 = 0 \quad (4)$$

$$\frac{d}{dr} (\rho \bar{v}_r^2 \bar{v}_E^2) + \rho \bar{v}_E^2 \frac{d\Phi}{dr} - \frac{\rho}{r} \bar{v}_E^4 + \frac{4\rho}{r} \bar{v}_r^2 \bar{v}_E^2 = 0 \quad (5)$$

[CHECK: FOR AN ISOTROPIC SYSTEM, IT TURNS OUT THAT $\bar{v}_r^2 \bar{v}_E^2 = \frac{2}{3} \bar{v}_r^4$ AND $\bar{v}_E^4 = \frac{8}{3} \bar{v}_r^4$

SO (4) $\Rightarrow \frac{d}{dr} (\rho \bar{v}_r^4) + 3\rho \bar{v}_r^2 \frac{d\Phi}{dr} = 0$

AND (5) $\Rightarrow \frac{d}{dr} (\rho \bar{v}_r^4) + \frac{3}{2} \rho \bar{v}_E^2 \frac{d\Phi}{dr} = 0$

I.E. AS $\bar{v}_E^2 = (\bar{v}_\theta^2 + \bar{v}_\phi^2) = 2\bar{v}_r^2$ FOR AN ISOTROPIC SYSTEM, THESE 2 EQUATIONS ARE DEGENERATE (AND RIGHT I THINK - SEE RESEARCH EXAM)

TO ELIMINATE $\bar{v}_r^2 \bar{v}_E^2$ BETWEEN (4) AND (5),

$$(4) \Rightarrow \rho \bar{v}_r^2 \bar{v}_E^2 = \frac{r}{3} \frac{d}{dr} (\rho \bar{v}_r^4) + \rho \bar{v}_r^2 r \frac{d\Phi}{dr} + \frac{2}{3} \rho \bar{v}_r^4$$

SUBSTITUTING INTO (5)

$$\Rightarrow \frac{d}{dr} \left[\frac{r}{3} \frac{d}{dr} (\rho \bar{v}_r^4) + \rho \bar{v}_r^2 r \frac{d\Phi}{dr} + \frac{2}{3} \rho \bar{v}_r^4 \right] + \rho \bar{v}_E^2 \frac{d\Phi}{dr} - \frac{\rho}{r} \bar{v}_E^4 + \frac{4}{3} \frac{d}{dr} (\rho \bar{v}_r^4) + 4\rho \bar{v}_r^2 \frac{d\Phi}{dr} + \frac{8}{3} \frac{1}{r} \rho \bar{v}_r^4 = 0$$

$$\Rightarrow \rho \bar{v}_E^4 = \frac{r^2}{3} \frac{d^2}{dr^2} (\rho \bar{v}_r^4) + \frac{7}{3} r \frac{d}{dr} (\rho \bar{v}_r^4) + \frac{8}{3} \rho \bar{v}_r^4$$

$$+ r^2 \frac{d\Phi}{dr} \frac{d}{dr} (\rho \bar{v}_r^2) + (5\rho \bar{v}_r^2 + \rho \bar{v}_E^2) r \frac{d\Phi}{dr} + \rho \bar{v}_r^2 r^2 \frac{d^2 \Phi}{dr^2}$$

BUT THE SECOND MOMENT EQUATION (HYDROSTATIC EQUILIBRIUM) YIELDS

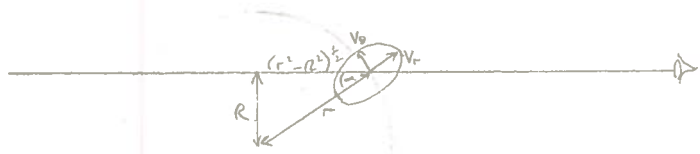
$$\frac{d}{dr} (\rho \bar{v}_r^2) + \frac{\rho}{r} [2\bar{v}_r^2 - \bar{v}_E^2] + \rho \frac{d\Phi}{dr} = 0$$

$$\Rightarrow \overline{\rho v_E^4} = \frac{r^2}{3} \frac{d^2}{dr^2} (\overline{\rho v_r^4}) + \frac{7}{3} r \frac{d}{dr} (\overline{\rho v_r^4}) + \frac{8}{3} \overline{\rho v_r^4} \quad (6)$$

$$+ \left\{ (3\overline{\rho v_r^2} + 2\overline{\rho v_E^2}) r \frac{d\Phi}{dr} - \rho r^2 \left(\frac{d\Phi}{dr} \right)^2 + \overline{\rho v_r^2} r^2 \frac{d^2\Phi}{dr^2} \right\}$$

THE QUANTITY WHICH CAN, IN PRINCIPLE, BE OBSERVED IN THIS SYSTEM IS THE PROJECTED LINE-OF-SIGHT 4TH VELOCITY MOMENT.

ONCE AGAIN FOLLOWING BINNEY AND TREMAINE (PP 207-8), THE PROJECTED 4TH MOMENT IS GIVEN BY



$$\mu(r) \overline{v_{los}^4} = 2 \int_R^\infty \overline{\rho(r) (v_r \cos \alpha - v_\theta \sin \alpha)^4} \frac{r dr}{(r^2 - R^2)^{1/2}}$$

$$= 2 \int_R^\infty \overline{\rho(r) [v_r^4 \cos^4 \alpha + 6v_r^2 v_\theta^2 \cos^2 \alpha \sin^2 \alpha + v_\theta^4 \sin^4 \alpha]} \times \frac{r dr}{(r^2 - R^2)^{1/2}}$$

BUT $\cos^2 \alpha = \frac{r^2 - R^2}{r^2}$; $\sin^2 \alpha = \frac{R^2}{r^2}$ AND $\overline{v_\theta^4} = \frac{3}{8} \overline{v_E^4}$; $\overline{v_r^2 v_\theta^2} = \frac{1}{2} \overline{v_r^2 v_E^2}$

$$\Rightarrow \mu(r) \overline{v_{los}^4} = 2 \int_R^\infty \left[\frac{(r^2 - R^2)^{3/2}}{r^3} \overline{\rho v_r^4} + \frac{3R^2 (r^2 - R^2)^{1/2}}{r^3} \overline{\rho v_r^2 v_E^2} + \frac{3}{8} \frac{R^4 (r^2 - R^2)^{-1/2}}{r^3} \overline{\rho v_E^4} \right] dr$$

SUBSTITUTING FOR $\overline{\rho v_r^2 v_E^2}$ FROM (4) AND $\overline{\rho v_E^4}$ FROM (6)

$$\Rightarrow \mu(r) \overline{v_{los}^4} = 2 \int_R^\infty \left\{ \frac{(r^2 - R^2)^{3/2}}{r^3} \overline{\rho v_r^4} + \frac{R^2 (r^2 - R^2)^{1/2}}{r^2} \left[\frac{d}{dr} (\overline{\rho v_r^4}) + 3\overline{\rho v_r^2} \frac{d\Phi}{dr} + \frac{2\rho}{r} \overline{v_r^4} \right] \right.$$

$$+ \frac{3}{8} \frac{R^4 (r^2 - R^2)^{-1/2}}{r^3} \left[\frac{1}{3} r^2 \frac{d^2}{dr^2} (\overline{\rho v_r^4}) + \frac{7}{3} r \frac{d}{dr} (\overline{\rho v_r^4}) + \frac{8}{3} \overline{\rho v_r^4} \right.$$

$$\left. \left. + (3\overline{\rho v_r^2} + 2\overline{\rho v_E^2}) r \frac{d\Phi}{dr} - \rho r^2 \left(\frac{d\Phi}{dr} \right)^2 + \overline{\rho v_r^2} r^2 \frac{d^2\Phi}{dr^2} \right] \right\} dr \quad (7)$$

RE-ARRANGING SO ALL THE "KNOWN QUANTITIES" ARE ON THE LHS

$$\begin{aligned} \Rightarrow \mu(r) \bar{v}_{\text{los}}^4 &= 2 \int_r^{\infty} \left(\frac{3R^2}{r^2} (r^2 - R^2)^{1/2} \rho \bar{v}_r^2 \frac{d\bar{v}_r}{dr} + \frac{3}{8} \frac{R^4}{r^3} (r^2 - R^2)^{-1/2} \left[3\rho \bar{v}_r^2 + 2\rho \bar{v}_r^2 r \frac{d\bar{v}_r}{dr} - \rho r^2 \left(\frac{d\bar{v}_r}{dr} \right)^2 + \rho \bar{v}_r^2 r^2 \frac{d^2 \bar{v}_r}{dr^2} \right] \right) dr \\ &= 2 \int_r^{\infty} \left(\frac{1}{8} \frac{R^4}{r} (r^2 - R^2)^{-1/2} \frac{d^2 (\rho \bar{v}_r^4)}{dr^2} + \frac{R^2}{r^2} (r^2 - R^2)^{1/2} (r^2 - \frac{1}{8} R^2) \frac{d(\rho \bar{v}_r^4)}{dr} + r (r^2 - R^2)^{-1/2} \rho \bar{v}_r^4 \right) dr \quad (8) \end{aligned}$$

SUBSTITUTING $x = r^2$, $x = R^2 \Rightarrow dr = \frac{1}{2} \frac{dx}{x^{1/2}}$; $\frac{d}{dr} = 2x^{1/2} \frac{d}{dx}$; $\frac{d^2}{dr^2} = 4x \frac{d^2}{dx^2} + 2 \frac{d}{dx}$

AND WRITING LHS = H(x) A MEASURABLE QUANTITY

$$\Rightarrow H(x) = \int_x^{\infty} (x-x)^{-1/2} \left[\frac{1}{2} x^2 \frac{d^2 \rho \bar{v}_r^4}{dx^2} + 2x \frac{d(\rho \bar{v}_r^4)}{dx} + \rho \bar{v}_r^4 \right] dx$$

INTEGRATING THE LAST TERM BY PARTS

$$\begin{aligned} &= \int_x^{\infty} (x-x)^{-1/2} \left[\frac{1}{2} x^2 \frac{d^2 \rho \bar{v}_r^4}{dx^2} + 2x \frac{d(\rho \bar{v}_r^4)}{dx} \right] dx + \left[\rho \bar{v}_r^4 \cdot 2(x-x)^{1/2} - \int 2(x-x)^{1/2} \cdot \frac{d(\rho \bar{v}_r^4)}{dx} dx \right]_x^{\infty} \\ &= \int_x^{\infty} (x-x)^{-1/2} \left[\frac{1}{2} x^2 \frac{d^2 \rho \bar{v}_r^4}{dx^2} + 2(2x-x) \frac{d(\rho \bar{v}_r^4)}{dx} \right] dx \end{aligned}$$

IF $\rho \bar{v}_r^4 \rightarrow 0$ FASTER THAN $x^{-1/2}$ AS $x \rightarrow \infty$
(TO BE CHECKED A POSTERIORI)

WRITING $y = \frac{d(\rho \bar{v}_r^4)}{dx}$

$$= \int_x^{\infty} (x-x)^{-1/2} \left[\frac{1}{2} x^2 \frac{dy}{dx} + 2(2x-x)y \right] dx$$

BY PARTS AGAIN

$$\begin{aligned} &= \int_x^{\infty} (x-x)^{-1/2} \left[\frac{1}{2} x^2 \frac{dy}{dx} + 2 \left[y \cdot \frac{2(-x+4x)}{3} (x-x)^{1/2} - \int \frac{2(-x+4x)}{3} (x-x)^{1/2} \frac{dy}{dx} dx \right] \right] dx \\ &= \int_x^{\infty} (x-x)^{-1/2} \left(\frac{4}{3} x^2 - \frac{20}{3} x + \frac{35}{6} x^2 \right) \frac{dy}{dx} dx \end{aligned}$$

IF $y \rightarrow 0$ FASTER THAN $x^{-3/2}$ AS $x \rightarrow \infty$
(I.E. SAME CRITERION AS ABOVE)

NOW TO ABEL INVERT THIS :

MULTIPLYING BY $(x-u)^{-1/2}$ AND INTEGRATING $\int_u^\infty dx$

$$\Rightarrow \int_u^\infty H(x) \frac{1}{(x-u)^{1/2}} dx = \int_{x=u}^\infty \int_{x=x}^\infty \frac{dy}{dy} \left(\frac{4x^2}{3} - \frac{20}{3} x^2 + \frac{35}{6} x^2 \right) \frac{1}{(x-x)^{1/2} (x-u)^{1/2}} dx dy$$

INTERCHANGING THE ORDER OF INTEGRATION

$$= \int_{y=u}^\infty \int_{x=u}^y \left(\frac{4}{3} x^2 - \frac{20}{3} x^2 + \frac{35}{6} x^2 \right) \frac{1}{(x-x)^{1/2} (x-u)^{1/2}} dx dy$$

$$= \int_{x=u}^\infty I(u, x) \frac{dy}{dy} dx$$

WHERE THE INNER INTEGRAL

$$I(u, x) = \int_{x=u}^x \left(\frac{35}{6} x^2 - \frac{20}{3} x^2 + \frac{4}{3} x^2 \right) \frac{1}{(-x^2 + (u+x)x - ux)^{1/2}} dx$$

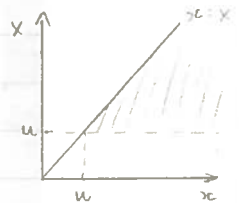
$$= \left[\frac{35}{6} \left(\frac{-2x - 3(u+x)}{4} \right) \frac{1}{(-x^2 + (u+x)x - ux)^{1/2}} + \frac{3(u+x)^2 - 4ux}{8} \left(\frac{dx}{(-x^2 + (u+x)x - ux)^{1/2}} \right) \right]$$

$$- \frac{20}{3} x \left(\frac{-1}{-1} \right) \frac{1}{(-x^2 + (u+x)x - ux)^{1/2}} + \frac{(u+x)}{2} \left(\frac{dx}{(-x^2 + (u+x)x - ux)^{1/2}} \right)$$

$$+ \frac{4}{3} x^2 \left(\frac{dx}{(-x^2 + (u+x)x - ux)^{1/2}} \right)$$

$$= \left(\frac{3}{16} x^2 - \frac{15}{8} ux + \frac{35}{16} u^2 \right) \left(\frac{dx}{(x-x)^{1/2} (x-u)^{1/2}} \right)$$

$$= \left(\frac{3}{16} x^2 - \frac{15}{8} ux + \frac{35}{16} u^2 \right) \cdot \pi$$



x > u

$$\Rightarrow \int_u^{\infty} \frac{H(x)}{(x-u)^{1/2}} dx = \int_{x=u}^{\infty} \frac{\pi}{16} (35u^2 - 30ux + 3x^2) \frac{dy}{dx} dx$$

DIFFERENTIATING W.R.T. u

$$\Rightarrow \frac{d}{du} \int_u^{\infty} \frac{H(x)}{(x-u)^{1/2}} dx = -\frac{\pi}{16} \cdot 8u^2 \frac{dy}{du} + \frac{\pi}{16} \int_{x=u}^{\infty} (70u - 30x) \frac{dy}{dx} dx$$

AND AGAIN

$$\begin{aligned} \Rightarrow \frac{d^2}{du^2} \int_u^{\infty} \frac{H(x)}{(x-u)^{1/2}} dx &= -\pi u \frac{dy}{du} - \frac{\pi}{2} u^2 \frac{d^2 y}{du^2} - \frac{\pi}{16} \cdot 40u \frac{dy}{du} + 70 \frac{\pi}{16} \int_{x=u}^{\infty} \frac{dy}{dx} dx \\ &= -\frac{\pi}{2} \left(u^2 \frac{d^2 y}{du^2} + 7u \frac{dy}{du} + \frac{35}{4} y \right) \end{aligned}$$

$$\Rightarrow u^2 \frac{d^2 y}{du^2} + 7u \frac{dy}{du} + \frac{35}{4} y = -\frac{2}{\pi} \frac{d^2}{du^2} \int_u^{\infty} \frac{H(x)}{(x-u)^{1/2}} dx$$

THIS IS AN EULER-TYPE EQUATION, SO LET $u = e^t \Rightarrow u \frac{d}{du} = \frac{d}{dt}$, $u^2 \frac{d^2}{du^2} = \frac{d^2}{dt^2} - \frac{d}{dt}$

$$\Rightarrow \frac{d^2 y}{dt^2} + 6 \frac{dy}{dt} + \frac{35}{4} y = -\frac{2}{\pi} e^{-2t} \left[\frac{d^2}{dt^2} - \frac{d}{dt} \right] \int_{e^t}^{\infty} \frac{H(x)}{(x-e^t)^{1/2}} dx$$

$$m^2 + 6m + \frac{35}{4} = 0$$

$$\Rightarrow m_{\pm} = \frac{-6 \pm \sqrt{36 - 35}}{2}$$

$$= -\frac{5}{2}, -\frac{7}{2}$$

WHICH HAS THE GENERAL SOLUTION

$$\begin{aligned} y &= c_1 e^{-5/2 t} + c_2 e^{-7/2 t} + \frac{e^{-5/2 t}}{1} \int e^{5/2 t'} \cdot -\frac{2}{\pi} e^{-2t'} \left[\frac{d^2}{dt'^2} - \frac{d}{dt'} \right] \int_{e^{t'}}^{\infty} \frac{H(x)}{(x-e^{t'})^{1/2}} dx dt' \\ &\quad + \frac{e^{-7/2 t}}{-1} \int e^{7/2 t'} \cdot -\frac{2}{\pi} e^{-2t'} \left[\frac{d^2}{dt'^2} - \frac{d}{dt'} \right] \int_{e^{t'}}^{\infty} \frac{H(x)}{(x-e^{t'})^{1/2}} dx dt' \end{aligned}$$

RE-SUBSTITUTING $u = e^t$ AND SUBSTITUTING $u' = e^{t'}$

$$\Rightarrow y = \cancel{c_1} u^{-5/2} + \cancel{c_2} u^{-7/2} - \frac{2}{\pi} u^{-5/2} \int u'^{5/2} \frac{d^2}{du'^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] \frac{du'}{u'} + \frac{2}{\pi} u^{-7/2} \int u'^{7/2} \frac{d^2}{du'^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] \frac{du'}{u'}$$

↑
ABSORBED INTO LOWER
LIMIT OF u' INTEGRATION

$$\Rightarrow y(u) = \frac{2}{\pi} \left\{ u^{-7/2} \int \frac{d^2}{du'^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] u'^{5/2} du' \right. \\ \left. - u^{-5/2} \int \frac{d^2}{du'^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] u'^{3/2} du' \right\}$$

$$\Rightarrow \frac{d}{dx} (\rho \sqrt{r^4}) = \frac{2}{\pi} \left\{ x^{-5/2} \int_x^{\infty} u'^{3/2} \frac{d^2}{du'^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] du' \right. \\ \left. - x^{-7/2} \int_x^{\infty} u'^{5/2} \frac{d^2}{du'^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] du' \right\} \quad x = r^2 \quad (9)$$

INTEGRATING THE FIRST INTEGRAL ON THE RHS BY PARTS

$$\Rightarrow \int_x^{\infty} u'^{3/2} \frac{d^2}{du'^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] du' \\ = \left[u'^{3/2} \cdot \frac{d}{du'} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] - \frac{3}{2} \int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right]_{u'=x}^{\infty}$$

ASSUMING THINGS ARE WELL-BEHAVED AT INFINITY, AND INTEGRATING BY PARTS AGAIN

$$= -x^{3/2} \frac{d}{dx} \left[\int_x^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx \right] - \frac{3}{2} \left[u'^{1/2} \cdot \int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx - \frac{1}{2} \int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right]_{u'=x}^{\infty} \\ = -x^{3/2} \frac{d}{dx} \left[\int_x^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx \right] + \frac{3}{2} x^{1/2} \int_x^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx + \frac{3}{4} \int_{u'=x}^{\infty} u'^{-1/2} \int_{x=u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx du'$$

BUT, REVERSING THE ORDER OF INTEGRATION ON THE LAST TERM,

$$\int_{u'=x}^{\infty} \int_{x=u'}^{\infty} u'^{-1/2} \frac{H(x)}{(x-u')^{1/2}} dx du' = \int_{x=x}^{\infty} H(x) \int_{u'=x}^x \frac{u'^{-1/2}}{(x-u')^{1/2}} du' dx \\ = \int_{x=x}^{\infty} H(x) \left[2 \tan^{-1} \sqrt{\frac{x}{x-1}} - 2 \tan^{-1} \sqrt{\frac{x}{x-1}} \right] dx$$

$$= \int_{x=xc}^{\infty} H(x) \cdot \left[2 \sin^{-1} \sqrt{\frac{x-xc}{x}} \right] dx$$

DOING SIMILAR THINGS WITH THE SECOND TERM,

$$\begin{aligned} \int_{xc}^{\infty} u^{5/2} \frac{d^2}{du^2} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] du' &= \left[u^{5/2} \cdot \frac{d}{du'} \left[\int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx \right] - \frac{5}{2} \int_{u'}^{\infty} \frac{H(x)}{(x-u')^{3/2}} dx \right] du' \Big|_{u'=xc}^{\infty} \\ &= -x^{5/2} \frac{d}{dx} \left[\int_{xc}^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx \right] - \frac{5}{2} \left[u^{3/2} \int_{u'}^{\infty} \frac{H(x)}{(x-u')^{1/2}} dx - \frac{3}{2} \int_{u'}^{\infty} \frac{H(x)}{(x-u')^{3/2}} dx du' \right] \Big|_{u'=xc}^{\infty} \\ &= -x^{5/2} \frac{d}{dx} \left[\int_{xc}^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx \right] + \frac{5}{2} x^{3/2} \int_{xc}^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx + \frac{15}{4} \int_{u'=xc}^{\infty} u^{1/2} \int_{x=u'}^{\infty} \frac{H(x)}{(x-u')^{3/2}} dx du' \end{aligned}$$

AGAIN REVERSING THE ORDER OF INTEGRATION ON THE LAST TERM,

$$\begin{aligned} \int_{u'=xc}^{\infty} \int_{x=u'}^{\infty} u^{1/2} \frac{H(x)}{(x-u')^{3/2}} dx du' &= \int_{x=xc}^{\infty} H(x) \int_{u'=xc}^x \frac{u^{1/2}}{(x-u')^{3/2}} du' dx \\ &= \int_{x=xc}^{\infty} H(x) \left[\sqrt{x(x-x)} + \frac{x}{2} \cdot \left[2 \sin^{-1} \sqrt{\frac{x-xc}{x}} \right] \right] dx \end{aligned}$$

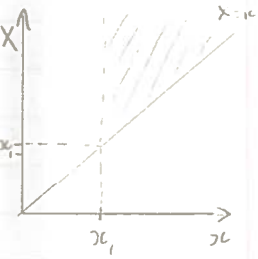
SUBSTITUTING ALL THIS BACK INTO (9)

$$\begin{aligned} \Rightarrow \frac{d}{dc} \rho \bar{v}_r^4 &= \frac{2}{\pi} \left\{ x^{-5/2} \left[-x^{3/2} \frac{d}{dc} \left(\int_{xc}^{\infty} \frac{H(x)}{(x-xc)^{1/2}} dx \right) + \frac{3}{2} x^{1/2} \int_{xc}^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx + \frac{3}{2} \int_{xc}^{\infty} \sin^{-1} \sqrt{\frac{x-xc}{x}} H(x) dx \right] \right. \\ &\quad \left. - x^{-3/2} \left[-x^{5/2} \frac{d}{dc} \left(\int_{xc}^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx \right) + \frac{5}{2} x^{3/2} \int_{xc}^{\infty} \frac{H(x)}{(x-x)^{1/2}} dx + \frac{15}{24} \int_{xc}^{\infty} \left(\sqrt{x(x-x)} \right. \right. \right. \\ &\quad \left. \left. \left. + x \sin^{-1} \sqrt{\frac{x-xc}{x}} \right) H(x) dx \right] \right\} \\ &= \frac{1}{\pi} \int_{x=xc}^{\infty} \left\{ -2x^{-2} \frac{H(x)}{(x-x)^{1/2}} - \frac{15}{2} x^{-3} (x-x)^{1/2} H(x) - \left(\frac{15}{2} x - 3x \right) x^{-3/2} \sin^{-1} \sqrt{\frac{x-xc}{x}} H(x) \right\} dx \end{aligned}$$

$$\Rightarrow \frac{d}{dc} \rho \bar{v}_r^4 = \frac{-1}{2\pi} x^{-7/2} \int_{x=xc}^{\infty} \left\{ (15x - 11x) x^{1/2} \frac{H(x)}{(x-x)^{1/2}} + (15x - 6x) \sin^{-1} \sqrt{\frac{x-xc}{x}} H(x) \right\} dx \quad (10)$$

SINCE $\rho \bar{V}_r^4 \rightarrow 0$ AS $r \rightarrow \infty$, WE CAN WRITE

$$\rho \bar{V}_r^4(x_1) = \frac{1}{2\pi} \int_{x=x_1}^{\infty} \int_{x=x_1}^{\infty} x^{-7/2} \left\{ (15x-11x_1) x^{1/2} (x-x_1)^{-1/2} + (15x-6x_1) \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) \right\} H(x) dx dx_1$$



REVERSING THE ORDER OF INTEGRATION (SIGH!)

$$= \frac{1}{2\pi} \int_{x=x_1}^{\infty} H(x) \int_{x=x_1}^x \left\{ 15x x^{-3} (x-x_1)^{-1/2} - 11x^{-2} (x-x_1)^{1/2} + 15x x^{-7/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) - 6x^{-5/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) \right\} dx dx$$

ATTACKING THE UNKNOWN INTEGRALS ONE AT A TIME

$a = -1$
 $b = x$

$$\int_{x=x_1}^x x^{-3} (x-x_1)^{-1/2} dx = \left[\frac{-3x-2x_1}{4x^2 x_1} (x-x_1)^{1/2} + \frac{3}{8x^2} \int x^{-1} (x-x_1)^{-1/2} dx \right]_{x_1}^x$$

$a = -1$
 $b = x$

$$\int_{x=x_1}^x x^{-2} (x-x_1)^{-1/2} dx = \left[-\frac{(x-x_1)^{1/2}}{x x_1} + \frac{1}{2x} \int x^{-1} (x-x_1)^{-1/2} dx \right]_{x_1}^x$$

$t = x^{1/2}$
 $\Rightarrow dx = 2t dt, x^{1/2} = t$

$$\int_{x=x_1}^x x^{-7/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) dx = 2 \int_{t=\sqrt{x_1}}^{\sqrt{x}} t^{-6} \cos^{-1}\left(\frac{t}{x^{1/2}}\right) dt$$

$$= 2 \left[\frac{t^{-5}}{-5} \cos^{-1}\left(\frac{t}{x^{1/2}}\right) - \frac{1}{5} \int t^{-5} (x-t^2)^{-1/2} dt \right]_{\sqrt{x_1}}^{\sqrt{x}}$$

$x = t^2$
 $\Rightarrow dt = \frac{1}{2} x^{-1/2} dx$

$$= \frac{2}{5} x_1^{-5/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) - \frac{1}{5} \int_{x_1}^x x^{-3} (x-x_1)^{-1/2} dx$$

$$\int_{x=x_1}^x x^{-5/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) dx = 2 \int_{t=\sqrt{x_1}}^{\sqrt{x}} t^{-4} \cos^{-1}\left(\frac{t}{x^{1/2}}\right) dt$$

$$= 2 \left[\frac{t^{-3}}{-3} \cos^{-1}\left(\frac{t}{x^{1/2}}\right) - \frac{1}{3} \int t^{-3} (x-t^2)^{-1/2} dt \right]_{\sqrt{x_1}}^{\sqrt{x}}$$

$$= \frac{2}{3} x_1^{-3/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) - \frac{1}{3} \int_{x_1}^x x^{-2} (x-x_1)^{-1/2} dx$$

SUBSTITUTING FOR THE TWO \cos^{-1} INTEGRALS, WE OBTAIN THE INTEGRAL OVER x

$$\begin{aligned}
 I_{\nu} &= 15x \int_{x_1}^x x^{-3}(x-x)^{-1/2} dx - 11 \int_{x_1}^x x^{-2}(x-x)^{-1/2} dx + 15x \left[\frac{2}{5} x_1^{-5/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) - \frac{1}{5} \int_{x_1}^x x^{-3}(x-x)^{-1/2} dx \right] \\
 &\quad - 6 \left[\frac{2}{3} x_1^{-3/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) - \frac{1}{3} \int_{x_1}^x x^{-2}(x-x)^{-1/2} dx \right] \\
 &= (6x - 4x_1) x_1^{-5/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) + 12x \int_{x_1}^x x^{-3}(x-x)^{-1/2} dx - 9 \int_{x_1}^x x^{-2}(x-x)^{-1/2} dx \\
 &= (6x - 4x_1) x_1^{-5/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) + \frac{3}{xx_1^2} (3x_1 + 2x)(x-x_1)^{1/2} + \frac{9}{2x} \int_{x_1}^x x^{-1}(x-x)^{-1/2} dx \\
 &\quad - \frac{9(x-x_1)^{1/2}}{2x} - \frac{9}{2x} \int_{x_1}^x x^{-1}(x-x)^{-1/2} dx
 \end{aligned}$$

THANKFULLY, THE DIVERGENT INTEGRALS CANCEL

$$\Rightarrow \rho \bar{\nu}_r^4(x) = \frac{1}{\pi} \cdot \frac{1}{x^2} \int_{x_1}^{\infty} \left\{ (3x - 2x) x^{-1/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) + 3(x-x)^{1/2} \right\} H(x) dx$$

FOR A WELL BEHAVED SYSTEM, $\rho \bar{\nu}_r^4$ IS FINITE AS $x \rightarrow 0$, AND THEREFORE THE INTEGRAL MUST VANISH UP TO $O(x^2)$ AS $x \rightarrow 0$. SERIES EXPANDING THE INTEGRAND YIELDS

$$\left\{ (3x - 2x) x^{-1/2} \cos^{-1}\left(\sqrt{\frac{x_1}{x}}\right) + 3(x-x)^{1/2} \right\} \approx \frac{3}{2} \pi x x^{-1/2} - \pi x^{1/2} + O(x^2)$$

SO $\rho \bar{\nu}_r^4$ WILL REMAIN FINITE AT THE ORIGIN ONLY IF

$$\int_{x=0}^{\infty} x H(x) dx = 0 \quad \text{AND} \quad \int_{x=0}^{\infty} H(x) dx = 0$$

WHICH LIVES THE 2 EXTRA MASS-TO-LIGHT RATIOS (SEE BELOW)

WE NOW NEED A CONVENIENT FORM FOR $H(x)$
 WITHOUT MESSING WITH HYDROSTATIC EQUILIBRIUM IN DERIVING (6), WE OBTAIN
 FOR EQUATION (7) THAT

$$\begin{aligned} \mu(r) \bar{V}_{\text{los}}^4 &= 2 \int_a^\infty \left\{ \frac{(r^2 - a^2)^{3/2}}{r^3} \rho \bar{V}_r^4 + \frac{R^2 (r^2 - a^2)^{1/2}}{r^2} \left[\frac{d(\rho \bar{V}_r^4)}{dr} + 3 \rho \bar{V}_r^2 \frac{d\Phi}{dr} + \frac{2e}{r} \bar{V}_r^4 \right] \right. \\ &\quad + \frac{3}{8} \frac{R^4}{r^3} (r^2 - a^2)^{-1/2} \left[\frac{1}{3} r^2 \frac{d^2(\rho \bar{V}_r^4)}{dr^2} + \frac{7}{3} r \frac{d(\rho \bar{V}_r^4)}{dr} + \frac{8}{3} \rho \bar{V}_r^4 \right. \\ &\quad \left. \left. + r^2 \frac{d\Phi}{dr} \frac{d(\rho \bar{V}_r^2)}{dr} + (5\rho \bar{V}_r^2 + \rho \bar{V}_e^2) r \frac{d\Phi}{dr} + \rho \bar{V}_r^2 r^2 \frac{d^2\Phi}{dr^2} \right] \right\} dr \end{aligned} \quad (7')$$

$$\begin{aligned} \Rightarrow \mu(r) \bar{V}_{\text{los}}^4 &= 2 \int_a^\infty \left\{ \frac{3R^2}{r^2} (r^2 - a^2)^{1/2} \rho \bar{V}_r^2 \frac{d\Phi}{dr} + \frac{3}{8} \frac{R^4}{r^3} (r^2 - a^2)^{-1/2} \left[r^2 \frac{d\Phi}{dr} \frac{d(\rho \bar{V}_r^2)}{dr} + (5\rho \bar{V}_r^2 + \rho \bar{V}_e^2) r \frac{d\Phi}{dr} \right. \right. \\ &\quad \left. \left. + \rho \bar{V}_r^2 r^2 \frac{d^2\Phi}{dr^2} \right] \right\} dr \\ &= 2 \int_a^\infty \left(\frac{1}{8} \frac{R^4}{r} (r^2 - a^2)^{-1/2} \frac{d^2 \rho \bar{V}_r^4}{dr^2} + \frac{R^2}{r^2} (r^2 - a^2)^{-1/2} (r^2 - \frac{1}{8} R^2) \frac{d(\rho \bar{V}_r^4)}{dr} + r (r^2 - a^2)^{-1/2} \rho \bar{V}_r^4 \right) dr \end{aligned} \quad (8')$$

SUBSTITUTING $xc = r^2$, $x = a^2$, WE CAN WRITE THE LHS OF (8') AS

$$H(x) = \mu \bar{V}_{\text{los}}^4(x) - 2 \int_x^\infty \left\{ \frac{3X(x-x)^{1/2}}{xc} \rho \bar{V}_r^2 F + \frac{3}{8} \frac{x^2}{x^{3/2}} (x-x)^{-1/2} \left[xF \cdot 2x^{1/2} \frac{d(\rho \bar{V}_r^2)}{dc} + (5\rho \bar{V}_r^2 + \rho \bar{V}_e^2) x^{1/2} F \right. \right. \\ \left. \left. + \rho \bar{V}_r^2 x \cdot 2x^{1/2} \frac{dF}{dc} \right] \right\} \frac{dc}{2x^{1/2}}$$

$$\text{WHERE } F(xc) = \frac{d\Phi}{dr} = \frac{GM(xc)}{xc} = G \gamma \frac{L(xc)}{xc}$$

$$\Rightarrow H(x) = \mu \bar{V}_{\text{los}}^4(x) - \int_x^\infty \frac{3}{8} \frac{x^2}{x^{3/2}} (xc - x)^{-1/2} \left[\left(\frac{8x}{x} - 3 \right) \rho \bar{V}_r^2 F + \rho \bar{V}_e^2 F + 2x \frac{d(\rho \bar{V}_r^2)}{dc} F + 2xc \rho \bar{V}_r^2 \frac{dF}{dc} \right] dc$$

NOTICE THAT THE FIRST TERM IN $H(x)$ DOES NOT DEPEND ON THE MASS-TO-LIGHT RATIO γ , WHEREAS THE SECOND TERM IS PROPORTIONAL TO γ VIA F . (THIS NEGLECTS THE FACT THAT THE MASS-TO-LIGHT RATIO WENT INTO CALCULATING $\overline{pV_r^3}$, $\overline{pV_r^2}$, BUT THIS CALCULATION IS ONLY BEING MADE TO CHECK FOR CONSISTENCY WITH THAT VALUE).

WE MAY THEREFORE WRITE

$$H(x) = \mu \overline{V_{\text{LOS}}^4}(x) - \gamma H_\gamma(x)$$

FOR A WELL BEHAVED SYSTEM (FOR EXAMPLE, ONE WHERE THE CENTRAL POTENTIAL IS FINITE, AND THE SYSTEM IS OF LIMITED EXTENT), THE 4TH MOMENT OF VELOCITY AT THE SYSTEM CENTER ($x=0$) WILL BE FINITE. WE THEREFORE CONSIDER THE EVALUATION OF (10) FOR SMALL x .

TAYLOR EXPANDING $(x-x)^{1/2}$ AND $\sin^{-1}(\sqrt{1-\frac{x}{x}}) = \cos^{-1}(\sqrt{\frac{x}{x}})$ IN THE INTEGRAND OF (10) ABOUT $x=0$ YIELDS A TERM

$$E = (15x - 11x) x^{1/2} \left(1 + \frac{1}{2} \frac{x}{x} + \frac{3}{8} \left(\frac{x}{x}\right)^2 + O(x^3) \right) + (15x - 6x) \left(\frac{\pi}{2} - \left(\frac{x}{x}\right)^{1/2} - \frac{1}{6} \left(\frac{x}{x}\right)^{3/2} - \frac{3}{40} \left(\frac{x}{x}\right)^{5/2} + O(x^3) \right)$$

MIRACULOUSLY, THIS SIMPLIFIES TO

$$E = \frac{15\pi x}{2} - 3\pi x + O(x^{3/2})$$

AND SO IT IS APPARENT THAT $\frac{d}{dx}(\overline{pV_r^3})$ WILL REMAIN FINITE AT $x=0$ ONLY IF

$$\int_{x=0}^{\infty} x H(x) dx = 0 \quad \text{AND} \quad \int_{x=0}^{\infty} H(x) dx = 0$$

OF THESE, THE SECOND IS MORE USEFUL, AS IT YIELDS THE MASS-TO-LIGHT RATIO

$$\gamma = \frac{\int_{x=0}^{\infty} \mu \overline{V_{\text{LOS}}^4} dx}{\int_{x=0}^{\infty} H_\gamma(x) dx} = \frac{\frac{1}{N_{\text{TOT}}} \langle \overline{V_{\text{LOS}}^4} \rangle}{\int_{x=0}^{\infty} H_\gamma(x) dx}$$

WHERE N_{TOT} IS THE TOTAL NUMBER OF PARTICLES IN THE SYSTEM AND $\langle \overline{V_{\text{LOS}}^4} \rangle$ IS THE GLOBAL MEAN 4TH MOMENT OF THE LINE-OF-SIGHT VELOCITY DISTRIBUTION, A COMPARATIVELY SIMPLE QUANTITY TO MEASURE. ACTUALLY, MAKES NO DIFFERENCE. USE BOTH!

NOW LET'S PRETTY UP THE TWO M/L INTEGRALS:

$$1) \int_{x=0}^{\infty} H_y(x) dx = \int_{x=0}^{\infty} \int_{x=X}^{\infty} \frac{3}{8} \frac{X^2}{x^{3/2}} (x-X)^{-1/2} \left[\left(\frac{8x}{X} - 3 \right) \rho \bar{v}_r^2 F + \rho \bar{v}_t^2 F + 2x \frac{d}{dc} (\rho \bar{v}_r^2 F) \right] dc dx$$

REVERSING THE ORDER OF INTEGRATION YET AGAIN

$$\begin{aligned} &= \int_{x=0}^{\infty} \int_{x=0}^x \frac{3}{8} \frac{X^2}{x^{3/2}} (x-X)^{-1/2} \left[\left(\frac{8x}{X} - 3 \right) \rho \bar{v}_r^2 F + \rho \bar{v}_t^2 F + 2x \frac{d}{dc} (\rho \bar{v}_r^2 F) \right] dx dx \\ &= \frac{3}{8} \int_{x=0}^{\infty} \frac{1}{x^{3/2}} \left[8x \rho \bar{v}_r^2 F \cdot \frac{4}{3} x^{3/2} + \left(-3 \rho \bar{v}_r^2 F + \rho \bar{v}_t^2 F + 2x \frac{d}{dc} (\rho \bar{v}_r^2 F) \right) \frac{16}{5} x^{5/2} \right] dx \\ &= \int_{x=0}^{\infty} \left[x \left(\frac{14}{5} \rho \bar{v}_r^2 F + \frac{2}{5} \rho \bar{v}_t^2 F \right) + \frac{4}{5} x^2 \frac{d}{dc} (\rho \bar{v}_r^2 F) \right] dx \end{aligned}$$

INTEGRATING THE LAST TERM BY PARTS

$$= \int_{x=0}^{\infty} \left[\frac{2x}{5} \left(7 \rho \bar{v}_r^2 F + \rho \bar{v}_t^2 F \right) + \frac{4}{5} \left[-2x \rho \bar{v}_r^2 F \right] \right] dx$$

$$\beta = 1 - \frac{1}{2} \frac{\bar{v}_t^2}{\bar{v}_r^2} \quad \Rightarrow \int_{x=0}^{\infty} H_y(x) dx = \int_{x=0}^{\infty} \frac{2}{5} x F(x) \left(3 \rho \bar{v}_r^2 + \rho \bar{v}_t^2 \right) dx = \int_{x=0}^{\infty} \rho \bar{v}_r^2 \left(1 - \frac{2}{5} \beta \right) F d(r^4)$$

$$2) \int_{x=0}^{\infty} x H_y(x) dx = \int_{x=0}^{\infty} x \int_{x=X}^{\infty} \frac{3}{8} \frac{X^2}{x^{3/2}} (x-X)^{-1/2} \left[\left(\frac{8x}{X} - 3 \right) \rho \bar{v}_r^2 F + \rho \bar{v}_t^2 F + 2x \frac{d}{dc} (\rho \bar{v}_r^2 F) \right] dx dx$$

REVERSING THE ORDER OF INTEGRATION

$$\begin{aligned} &= \int_{x=0}^{\infty} \int_{x=0}^x \frac{3}{8} \frac{X^3}{x^{3/2}} (x-X)^{-1/2} \left[\left(\frac{8x}{X} - 3 \right) \rho \bar{v}_r^2 F + \rho \bar{v}_t^2 F + 2x \frac{d}{dc} (\rho \bar{v}_r^2 F) \right] dx dx \\ &= \frac{3}{8} \int_{x=0}^{\infty} \frac{1}{x^{3/2}} \left[8x \rho \bar{v}_r^2 F \cdot \frac{16}{15} x^{5/2} + \left(-3 \rho \bar{v}_r^2 F + \rho \bar{v}_t^2 F + 2x \frac{d}{dc} (\rho \bar{v}_r^2 F) \right) \frac{32}{35} x^{7/2} \right] dx \end{aligned}$$

$$\Rightarrow \int_{x=0}^{\infty} x H_Y(x) dx = \int_{x=0}^{\infty} \left[x^2 \left(\frac{76}{35} \rho \bar{v}_r^2 F + \frac{12}{35} \rho \bar{v}_t^2 F \right) + \frac{24}{35} x^3 \frac{d}{dx} (\rho \bar{v}_r^2 F) \right] dx$$

INTEGRATING THE LAST TERM BY PARTS

$$= \int_{x=0}^{\infty} \left[\frac{4}{35} x^2 (19 \rho \bar{v}_r^2 F + 3 \rho \bar{v}_t^2 F) - \frac{72}{35} x^2 \rho \bar{v}_r^2 F \right] dx$$

$$\Rightarrow \int_{x=0}^{\infty} x H_Y(x) dx = \int_{x=0}^{\infty} \frac{4}{35} x^2 F(x) (\rho \bar{v}_r^2 + 3 \rho \bar{v}_t^2) dx = \int_{x=0}^{\infty} \frac{4}{15} \rho \bar{v}_r^2 (1 - \frac{6}{7} \beta) F d(r^6)$$

TO SUMMARIZE, THE ESTIMATIONS FOR G*M/L ARE:

$$\gamma_0 = \frac{3}{2} \frac{\int_0^{\infty} R \mu \bar{v}_{LOS}^2 dR}{\int_0^{\infty} r v L dr}$$

$$\gamma_1 = \frac{1}{2} \frac{\int_0^{\infty} R \mu \bar{v}_{LOS}^4 dR}{\int_0^{\infty} r v \bar{v}_r^2 (1 - \frac{2}{5} \beta) L dr}$$

$$\gamma_2 = \frac{5}{4} \frac{\int_0^{\infty} R^3 \mu \bar{v}_{LOS}^4 dR}{\int_0^{\infty} r^3 v \bar{v}_r^2 (1 - \frac{6}{7} \beta) L dr}$$